### **Dibosons at the LHC** Symmetry Restored?

Joachim Kopp Invisibles Webinar I November 24<sup>th</sup> 2015

based on arXiv:1507.00013 (with J. Brehmer, J. Hewett, T. Rizzo, J. Tattersall)





### Diboson Anomalies at the LHC

A Global Fit to the Data

### Theoretical Explanations on the Market

### Left-Right Symmetric Models

### **M** Conclusions





2

## Diboson Anomalies at the LHC





3

#### A New Resonance at 1.8–2 TeV?



- 3.4 $\sigma$  Resonance in  $pp \rightarrow (V \rightarrow jj) + (V \rightarrow jj)$  at 1.8–2 TeV
- Poor discrimination between W and Z





#### **Boosted WW, ZZ and WZ**













6 JG









## A Global Fit to the Data





#### **Cross Section Fit: Diboson Final States**

- Assume heavy resonance decaying to WW, ZZ or WZ
- Combine all searches for these final states:
  - ATLAS  $VV \rightarrow JJ$  (3.4 $\sigma$ ) 1506.00962
  - CMS  $VV \rightarrow JJ$  (~1 $\sigma$ ) 1405.1994
  - ATLAS  $WW/WZ \rightarrow \ell\nu + jets$  1503.04677
  - CMS  $VV \rightarrow \ell \nu / \ell^+ \ell^- + \text{jets}$  1405.3447
  - ATLAS  $VV \rightarrow \ell^+ \ell^- + jets$  1409.6190
- Input Data
  - # of events + predicted BG in mass window [1.7 ... 2.0 TeV]
  - Acceptance, Efficiency, Systematics where available
  - where not: tune to reproduce experimental limits



8

#### **Cross Section Fit: Diboson Final States**



 $\sigma(pp \rightarrow X) \times BR(X \rightarrow ZZ) \text{ [fb]}$ 



JG

#### **Cross Section Fit: V+H Final States**

- Assume heavy resonance decaying to WH or ZH
- Combine all searches for these final states:
  - ATLAS  $VH \rightarrow b\overline{b} + \ell\ell/\ell\nu/\nu\nu$ 1503.08089
  - CMS  $WH \rightarrow b\bar{b} + \ell\nu$  (~2.1 $\sigma$ ) CMS-PAS-EXO-14-010
  - CMS  $VH \rightarrow \tau^+\tau^- + jj$ 1502.04994
  - CMS  $VH \rightarrow JJ$ 1506.1443







#### **Cross Section Fit: Dijet and t+b Resonances**

- Assume heavy resonance decaying to jj or tb
- Combine all searches for these final states:
  - ATLAS dijet (~2σ)
     1407.1376
  - CMS dijets (~2σ)
     1501.04198
  - ATLAS  $tb \rightarrow hadrons$  1408.0886
  - CMS  $tb \rightarrow \ell \nu bb$ 1402.2176





- ZZ or WZ resonance favored over WW
  - σ ~ 5 fb
  - New boson, preferably charged
- Possible resonance in WH (but not ZH)
  - σ ~ 5 fb
- Dijet resonance (constraints form tb channel)
  - σ ~ 100 fb
  - Strong coupling to quarks
- No resonance observed in (semi)leptonic final states
  - suppressed couplings to leptons





## Explanations





#### **Explanations in the Literature**

- A new composite sector
- An enlarged scalar sector
- Extended gauge groups
- Others



Fukano et al., 1506.03751: Thamm Torre Wulzer, 1506.08688: Carmona Delgadod Quiros Santiago, 1507.01914: Chiang Fukuda Harigaya Ibe Yanagida, 1507.02483: Cacciapaglia Deandrea Hashimoto, 1507.03098: Sanz, 1507.03553: Bian Liu Shu, 1507.06018: Low Tesi Wang, 1507.07557: Dobadu Guo Llanes-Estrada, 1508.03544:

#### techni-rho

composite Higgs sector

composite Higgs sector

composite scalar

composite pseudoscalar, WZW term

glueball in new composite sector

composite Higgs sector

composite Higgs sector

composite Higgs sector





Hisano Nagata Omura, 1506.03931: Cheung Keung Tseng Yuan, 1506.06064: Dobrescu Liu, 1506.06736 / 1507.01923: Gao Ghosh Sinha Yu, 1506.07511: Brehmer Hewett JK Rizzo Tattersall, 1507.00013: Abe Nagai Okawa Tanabashi, 1507.0115: Cao Yan Zhang, 1507.00268: Heeck Patra, 1507.01584: Abe Kitahara Nojiri, 1507.01681: Anchordoqui et al., 1507.05299: Dhuria Hati Sarkar, 1507.08297: Dev Mohapatra, 1508.02277: Coloma Dobrescu Lopez-Pavon, 1508.04129: Deppisch *et al.*, 1508.05940:

Leptophobic Z' or W' Z' or  $W_{R}$ W' in LR model W' in LR model W' in LR model KK-W in 3-site moose model Various 221 and 331 models W' in LR model strongly interacting SU(2)' stringy U(1)' string-inspired U(1)' W' in LR model W' in LR model W' in LR model



Bian Liu Shu Zhang, 1509.02787:
Bandyopadhyay *et al.*, 1509.03232:
Avasthi Dev Mitra, 1509.05387:
Li Maxin Mayes Nanopoulos, 1509.06821:
Ko Nomura, 1510.07872:
Collins Ng, 1510.08083:
Dobrescu Fox, 1511.02148:
Wang Sage Steele Mann, 1511.02531:
Appelquist Bai Ingoldby Piai, 1511.05473:

W'

W' in LR model / SO(10) context
W' in LR model
stringy leptophobic U(1)'
W' in LR model
W' in SUSY LR model
W' in LR model
U(1)' + classical conformal symmetry
SU(2)<sub>L</sub> x SU(2)<sub>L</sub> x SU(2)<sub>R</sub> x SU(2)<sub>R'</sub> x U(1)<sub>Y</sub>





Chen Nomura, 1507.04431: Omura Tobe Tsumura, 1507.05028: Chao, 1507.05310: Chen Nomura, 1509.02039: Sierra *et al.*, 1510.03437: Alves Camargo Dias, 1511.04449:

2HDM

2HDM (H<sup>0</sup> with large up-Yukawa)

2HDM

Higgs singlet (+ vector-like fermions) 2HDM

Heavy Higgs (+ vector-like fermions)





Aguilar-Saavedra, 1506.06739: Cacciapaglia Franssen, 1507.00900: Fukano Matsuzaki Yamawaki, 1507.03428: Kim Kong Lee Park, 1507.06312: Liew Shirai, 1507.08273: Arnan Espriu Mescia, 1508.00174: Goncalves Krauss Spannowsky, 1508.04162: Fichet Gersdorff, 1508.04814: Petersson Torre, 1508.05632: Arbuzov Zaitsev, 1510.02312: Allanach Dev Sakurai, 1511.01483: Sajjad, 1511.02244: Bhattacherjee *et al.*, 1511.02797:

triboson final state (X -> W + Z + Y) effective Lagrangian with heavy vector effective Lagrangian with heavy vector EFT analysis RS graviton Model-independent unitarity arguments

Model-indep. results on jet substructure

EFT for spin-0 and spin-2 resonances

sgoldstino

Heavy scalar with dim-5 couplings

RPV SUSY with  $m_{ au} \sim m_W, m_Z$ 

EFT analysis

triboson final state (X  $\rightarrow$  W + Z + Y)



19

#### **Composite Higgs Models**

- SM fermions mix with fermions of heavy composite sector
- Symmetry breaking e.g.  $SO(5) \rightarrow SO(4)$ 
  - Higgs as pseudo-Goldstone boson
- Heavy resonances at compositeness scale
  - One of them could explain the diboson anomaly
  - Typically described using a phenomenological Lagrangian

$$\begin{split} \mathcal{L} \supset &-\frac{1}{4} D_{[\mu} V_{\nu]}^{a} D^{[\mu} V^{\nu] a} + \frac{m_{V}^{2}}{2} V_{\mu}^{a} V^{\mu a} \\ &+ i g_{V} c_{H} V_{\mu}^{a} H^{\dagger} \tau^{a} \overleftrightarrow{D}^{\mu} H + \frac{g^{2}}{g_{V}} c_{F} V_{\mu}^{a} J_{F}^{\mu a} \\ &+ \frac{g_{V}}{2} c_{VVV} \epsilon_{abc} V_{\mu}^{a} V_{\nu}^{b} D^{[\mu} V^{\nu] c} + g_{V}^{2} c_{VVHH} V_{\mu}^{a} V^{\mu a} H^{\dagger} H \\ &- \frac{g}{2} c_{VVW} \epsilon_{abc} W^{\mu \nu a} V_{\mu}^{b} V_{\nu}^{c} \end{split}$$

Thamm Torre Wulzer, 1506.08688



Carmona Delgadod Quiros Santiago,1507.01914



#### Fit of a Composite Vector Resonance



Carmona Delgadod Quiros Santiago, 1507.01914





# Left-Right Symmetry























 $W_L, Z_L = W_R, Z_R$ Higgs sector  $SU(2)_L \times SU(2)_R \times U(1)_{B-L}$  $SU(2)_L \times U(1)_Y$  $W_L, Z_L$  mass  $W_L-W_R, Z_L-Z_R$  mixing  $\Phi$  (2, 2, 0) physical states W, W' $U(1)_{\rm em}$ 

> $\Delta_L \sim (3, 1, 2)$ (zero or tiny vev)



23





#### W / W' Phenomenology in the LR Model

• W boson mass matrix

$$\mathcal{M}_{\mathcal{W}}^2 = \begin{pmatrix} m_{W_L}^2 & \beta_w m_{W_L}^2 \\ \beta_w m_{W_L}^2 & m_{W_R}^2 \end{pmatrix}$$

Physical masses

 $m_1 \sim m_{W_L} \qquad m_2 \sim m_{W_R}$ 

• Mixing angle

$$\tan 2\phi_w = \frac{-2\beta_w m_W^2}{m_{W_R}^2 - m_W^2}$$

•  $Z_R$  mass

$$\frac{m_{Z_R}^2}{m_{W_R}^2} = \frac{2\kappa^2(1 - \sin^2\theta_w)}{\kappa^2(1 - \sin^2\theta_w) - \sin^2\theta_w} > 1$$





#### W / W' Phenomenology in the LR Model



#### W' and Z' Couplings to Fermions

$$Q = T_{3L} + T_{3R} + (B - L)/2$$

$$\mathcal{O}_{W'} \simeq rac{g_R}{\sqrt{2}}$$
 (neglecting  $W - W'$  mixing)

$$W_R \sim q'$$

• Z' couplings

$$\mathcal{O}_{Z_R} = \frac{g_R}{\cos \alpha} \left[ T_{3R} + Y \sin^2 \alpha \right] \quad \text{with} \quad \tan \alpha \equiv \frac{g_{B-L}}{g_R}$$



#### W'-W-Z Coupling

• Covariant derivative:

$$\mathcal{D} = \partial - ieQA - \frac{i}{\sqrt{2}}g_L T_L^{\pm} \cdot W^{\mp} + (L \to R)$$
$$-i\frac{g_L}{\cos\theta_w}(T_{3L} - \sin^2\theta_w Q)Z - i\mathcal{O}_{Z_R} Z_R$$

- Two contributions to the W' W Z coupling:
  - $\mathcal{D}$  acting on  $W_L^+ W_L^-$ :  $\sin \phi_w \cos \phi_w \times \frac{g_L}{\cos \theta_w} (1 \sin^2 \theta_w)$



 $\approx W_R \sim ($ 

#### Fitting the LR Model







#### Fitting the LR Model





#### The CMS eejj anomaly

- Added bonus: explanation of the CMS eejj anomaly (2.8 $\sigma$ ) via  $W_R \rightarrow eN \rightarrow ee + (W_R^* \rightarrow jj)$
- Potential problems
  - Anomaly only in e<sup>+</sup>e<sup>-</sup>, not e<sup>±</sup>e<sup>±</sup> possible remedy: Dirac N
  - No peak in m<sub>ejj</sub> distribution
  - No signal in µµjj

Dobrescu Liu, 1506.06736 Dobrescu Liu, 1507.01923 Deppisch *et al.*, 1508.05940 Dobrescu Fox, 1511.02148





#### **Prospects at 13 TeV**

- W<sub>R</sub> production cross section ~6 times larger than at 8 TeV
- Best fit W<sub>R</sub> can be excluded with
  - ▶ 5 fb<sup>-1</sup> in jj
  - 10 fb<sup>-1</sup> in tb
  - 15 fb<sup>-1</sup> in WZ, WH
- Predictions for jj at best fit
  - $\sigma \times BR \times A \sim 300 \text{ fb}$





## Connections to Dark Matter?





#### W'-mediated coupling to SM fields

- DM candidate  $\chi^0$  (e.g.  $N_3$ )
- $SU(2)_R$  partner  $\chi^{\pm}$  (e.g.  $\tau_R$ )
- Requirements



- $m_{N_3} < m_{ au}$  to prevent DM decay through off-shell  $W_R$
- No mixing between  $N_3$  and  $N_{1,2}$  to forbid fast decay to  $e/\mu + W_R^*$
- Entropy production after DM freeze-out  $(\langle \sigma v \rangle (N_3 N_3 \rightarrow \tau^+ \tau^-)$  too small)





#### A new SU(2)<sub>R</sub> doublet

- DM candidate  $\chi^0$
- In new SU(2)<sub>R</sub> doublet with charged partner  $\chi^{\pm}$
- Features
  - Thermal freeze-out works if  $m_{\chi^\pm} \sim m_{\chi^0}$  (coannihilation)
  - No direct or indirect signals
  - $BR(W_R \rightarrow \text{visible})$  reduced => larger  $\kappa$  allowed
  - Mixing between  $\chi^0$  and  $N_j$  must vanish to forbid DM decay





#### **Connections to Dark Matter?**

#### Minimal Left-Right Symmetric Dark Matter

- New SU(2)<sub>R</sub> triplet or quintuplet
- Mass splitting between  $\chi^0$ ,  $\chi^{\pm}$ ,  $\chi^{\pm\pm}$  from radiative corrections



Heeck Patra, 1507.01584 Ko, Nomura, 1510.07872



#### • DM in SUSY LR Models

 SUSY versions of the SU(2)<sub>L</sub> x SU(2)<sub>R</sub> model admit the usual SUSY DM candidates





#### Z- and Z<sub>R</sub>-mediated DM interactions

- New SU(2)<sub>R</sub> multiplet  $(\chi^0, \chi^{\pm})$
- Extra  $\mathbb{Z}_2$  symmetry forbids  $\chi^0 N_j$  mixing
- Features
  - $\chi^0$  has Majorana mass => DM—nucleus scattering spin-dependent

$$\sigma_{\chi N} = \frac{3m_N^2 m_{\chi}^2}{\pi (m_N + m_{\chi})^2} \left[ \sum_{q=u,d,s} \Delta_{Nq} \left( \frac{g_{qA} g_{\chi A}}{M_Z^2} + \frac{g_{qA}' g_{\chi A}'}{M_Z'^2} \right) \right]^2$$

 $\blacktriangleright$  Annihilation through Z and  $Z_{R}$ 

$$\langle \sigma v \rangle (\chi \chi \to f \bar{f}) \simeq \frac{n_c v_{\rm rel}^2}{6\pi} \frac{g_{\chi A}^{\prime 2} (g_{fV}^{\prime 2} + g_{fA}^{\prime 2}) m_{\chi}^2}{(4m_{\chi}^2 - m_{Z_R}^2)^2 + m_{Z_R}^2 \Gamma_{Z_R}^2}$$







#### Z- and Z<sub>R</sub>-mediated DM interactions

- New SU(2)<sub>R</sub> multiplet  $(\chi^0, \chi^{\pm})$
- Extra  $\mathbb{Z}_2$  symmetry forbids  $\chi^0 N_j$  mixing
- Features
  - $\chi^0$  has Majorana mass => DM—nucleus scattering spin-dependent

$$\sigma_{\chi N} = \frac{3m_N^2 m_{\chi}^2}{\pi (m_N + m_{\chi})^2} \left[ \sum_{q=u,d,s} \Delta_{Nq} \left( \frac{g_{qA} g_{\chi A}}{M_Z^2} + \frac{g_{qA}' g_{\chi A}'}{M_Z'^2} \right) \right]^2$$

 $\blacktriangleright$  Annihilation through Z and  $Z_{R}$ 

$$\langle \sigma v \rangle (\chi \chi \to f \bar{f}) \simeq \frac{n_c v_{\rm rel}^2}{6\pi} \frac{g_{\chi A}^{\prime 2} (g_{fV}^{\prime 2} + g_{fA}^{\prime 2}) m_{\chi}^2}{(4m_{\chi}^2 - m_{Z_R}^2)^2 + m_{Z_R}^2 \Gamma_{Z_R}^2}$$







#### Z- and Z<sub>R</sub>-mediated DM interactions

- New SU(2)<sub>R</sub> multiplet  $(\chi^0, \chi^{\pm})$
- Extra  $\mathbb{Z}_2$  symmetry forbids  $\chi^0 N_j$  mixing
- Features
  - $\chi^0$  has Majorana mass => DM—nucleus scattering spin-dependent

$$\sigma_{\chi N} = \frac{3m_N^2 m_{\chi}^2}{\pi (m_N + m_{\chi})^2} \left[ \sum_{q=u,d,s} \Delta_{Nq} \left( \frac{g_{qA} g_{\chi A}}{M_Z^2} + \frac{g_{qA}' g_{\chi A}'}{M_Z'^2} \right) \right]^2$$

 $\blacktriangleright$  Annihilation through Z and  $Z_{R}$ 

$$\langle \sigma v \rangle (\chi \chi \to f \bar{f}) \simeq \frac{n_c v_{\rm rel}^2}{6\pi} \frac{g_{\chi A}^{\prime 2} (g_{fV}^{\prime 2} + g_{fA}^{\prime 2}) m_{\chi}^2}{(4m_{\chi}^2 - m_{Z_R}^2)^2 + m_{Z_R}^2 \Gamma_{Z_R}^2}$$



## Conclusions





#### Conclusions

#### $\mathbf{M}$ An intriguing accumulation of anomalies at $\sim 1.8 - 2$ TeV

- Signs of a composite sector?
- An extended Higgs sector?
- A new gauge boson?

#### Left-right symmetric interpretation

- Established and well-motivated model
- Good fit to the data
- Possible connections to the dark sector





## Thank You!





### Bonus Slides





#### **Cross Section Fit: Dijet and t+b Resonances**

- Assume heavy resonance decaying to jj or tb
- Combine all searches for these final states:
  - ATLAS dijet (~2σ)
     1407.1376
  - CMS dijets (~2σ)
     1501.04198
  - ATLAS  $tb \rightarrow hadrons$  1408.0886
  - ATLAS  $tb \rightarrow \ell \nu bb$ (systematically too strong by ~1.8 $\sigma$ ) 1410.4103
  - CMS  $tb \rightarrow \ell \nu bb$ 1402.2176





#### Fitting the LR Model



